# Spherically Symmetric Quantities in Biometric Geometry 

S. S. Charjan<br>Department of Mathematics, K. Z. S. Science College, Bramhani Kalmeshwar, Dist. Nagpur, India

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#### Abstract

A geometry based on two metrics at each point of space is proposed purely from mathematical point of view. Defining the bimetrically Spherically Symmetric Scalars, Vectors and Tensors of order -2 are deduced in bimetric geometry.


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## INTRODUCTION

As far as our knowledge goes there are some attempts of introducing two metric at each point of space; Rosen (1940), Eisenhart (1960) and Logunov and Mestvirishvili (1989). The geometry to that effect was utilized by Rosen (1940)in proposing theory Known as bimetric theory of gravitation or bimetric theory of relativity of Rosen, which differs from the general theory of relativity of Einstein $(1912,1915)$ and which was later explored by many workers [ Falik and Rosen (1980,1981), De Liebscer (1975), Yilmaz (1977), Goldman and Rosen (1977), Rosen and Rosen (1977), Israelit and Rosen (1983),Israelit (1979,1981), Rosen (1973,1974,1975,1979,1980,1983)]. This theory has some advantages over the general relativity: the quantities such as Chriostffel symbols and others becomes tensors which otherwise in Riemannian geometry they are not. Also on physical grounds there is no space for black holes. Our approach in constructing the bimetric geometry is rather different- it is purely mathematical. However, we shall be studying the impact of Rosen's bimetric theory with regard to our geometric ideas and the relevance of various bimetric quantities constructed in due course.

The concept of spherical symmetry (s. s.) in general relativity have been adopted from that of Euclidean geometry .Takeno (1966) has put the concept of spherical symmetry purely on a mathematical stand point. In his work the behavior of the spherically symmetric quantities is relegated to vanishing of Lie derivative and in consequence of which the forms of physical quantities (scalar, vector, tensor of rank 2 ) have been investigated.

In this paper, the attempt have been made to review this work of Takeno from the geometry of a bimetric theory of relativity proposed by Rosen (1940), and the related geometry contemplated by Karade and Nahatkar (2001).

## DEFINITION.

A geometric quantity $\Omega$ is spherically symmetric with regard to bimetric geometry (or a quantity $\Omega$ is said to be bimetrically spherically symmetric) if

$$
\begin{equation*}
\nabla_{m(\xi)} \Omega=0 \tag{1}
\end{equation*}
$$

for $\xi=\xi^{a i},(a=1,2,3)$, where $\nabla_{m(\xi)}$ is the bimetric covariant derivative with respect to $\xi$ and is given by $\nabla_{m(\xi)}=\left(g_{m}-f_{m}\right)_{(\xi)}, g_{m}=g$-covariant derivative with respect to metric $g_{i j}$ and $f_{m}=f$-covariant derivative with respect to metric $f_{i j}$. In this definition, $\xi^{a i}$ are the Killing vectors given by
$\xi^{1 i}=(0, \sin \phi, \cot \theta \cos \phi, 0), \quad \xi^{2 i}=(0,-\cos \phi, \cot \theta \sin \phi, 0), \xi^{3 i}=(0,0,-1,0)$.

## SPHERICALLY SYMMETRIC QUANTITIES OF THE SPHERICALLY SYMMETRIC SPACE-TIME V4.

## Scalar.

Following (1), the scalar $A$ is bimetrically spherically symmetric if $\nabla_{m(\xi)} A=\left(g_{m}-f_{m}\right)_{(\xi)} A=0$

$$
\begin{align*}
& \nabla_{m(\xi)} A=\xi^{m} \frac{\partial A}{\partial x^{m}}-\xi^{m} \frac{\partial A}{\partial x^{m}}=0 \\
& \nabla_{m(\xi)} A=\xi^{m}\left(\frac{\partial A}{\partial x^{m}}-\frac{\partial A}{\partial x^{m}}\right)=0 \tag{3}
\end{align*}
$$

and it is holds trivially, for any value of $A$. Hence the bimetric Spherically Symmetric scalar $A$ can be taken as a function of $r$ and $t$ to match with Takeno (1966).

## Covariant vector.

Following (1), the covariant vector $A_{i}$ is bimetrically spherically symmetric if

$$
\begin{align*}
& \nabla_{m(\xi)} A_{i}=\left(g_{m}-f_{m}\right)_{(\xi)} A_{i}=0 \\
& \nabla_{m(\xi)} A=\xi^{m} \frac{\partial A_{i}}{\partial x^{m}}+A_{m} \frac{\partial \xi^{m}}{\partial x^{i}}-\xi^{m} \frac{\partial A_{i}}{\partial x^{m}}=0 \\
& \nabla_{m(\xi)} A=A_{m} \frac{\partial \xi^{m}}{\partial x^{i}}=0 .  \tag{4}\\
& \text { ie } A_{m} \frac{\partial \xi^{a m}}{\partial x^{i}}=0 \tag{5}
\end{align*}
$$

which on simplification gives

$$
\begin{aligned}
& A_{1} \cdot 0+A_{2} \cdot 0+A_{3} \cdot 0+A_{4} \cdot 0=0 \\
& -A_{3} \cos c e^{2} \theta \cos \phi=0 \quad \Rightarrow A_{3}=0
\end{aligned}
$$

$$
A_{2} \cos \phi-A_{3} \cot \theta \sin \phi=0 \Rightarrow A_{2}=0
$$

Thus the bimetrically spherically symmetric vector $A_{i}$ have the form $A_{i}=\left(A_{1}, 0,0, A_{4}\right)$ in which $A_{1}$ and $A_{4}$ are the functions of $r, \theta, \phi, t$. If one takes $A_{1}$ and $A_{4}$ independent of $\theta$ and $\phi$, the result of Takeno (1966) is reclaimed.

## Contravarient Vector $A^{\mathrm{i}}$

Following (1), the contravarient vector $A^{i}$ is bimetrically spherically symmetric if

$$
\begin{align*}
\nabla_{m(\xi)} A^{i} & =\left(g_{m}-f_{m}\right)_{(\xi)} A^{i}=0 \\
\nabla_{m(\xi)} A & =\xi^{m} \frac{\partial A^{i}}{\partial x^{m}}-A^{m} \frac{\partial \xi^{i}}{\partial x^{m}}-\xi^{m} \frac{\partial A^{i}}{\partial x^{m}}=0 \\
\nabla_{m(\xi)} A & =-A^{m} \frac{\partial \xi^{i}}{\partial x^{m}}=0 \\
\text { ie } \quad \nabla_{m(\xi)} A_{i} & =A^{m} \frac{\partial \xi^{a i}}{\partial x^{m}}=0 \\
A^{m} \frac{\partial \xi^{a i}}{\partial x^{m}} & =0 \tag{6}
\end{align*}
$$

Which on simplification gives

$$
\begin{aligned}
& A^{1} \cdot 0+A^{2} .0+A^{3} .0+A^{4} .0=0 \\
& A^{3} \cos \phi=0 \Rightarrow A^{3}=0 \\
& A^{2}\left(-\cos c e^{2} \theta \cos \phi\right)=0 \Rightarrow A^{2}=0
\end{aligned}
$$

Thus the bimetrically spherically symmetric vector $A^{i}$ have the form $A^{i}=\left(A^{1}, 0,0, A^{4}\right)$ in which $A^{1}$ and $A^{4}$ are the functions of $r, \theta, \phi, t$. If one takes $A^{1}$ and $A^{4}$ independent of $\theta$ and $\phi$, the result of Takeno (1966) is reclaimed.

## Covariant tensor of order-2.

In view of definition (1), covariant tensor $A_{i j}$ of order- 2 is bimetrically spherically symmetric if

$$
\begin{gather*}
\nabla_{m(\xi)} A_{i j}=\left(g_{m}-f_{m}\right)_{(\xi)} \quad A_{i j}=0 \\
\nabla_{m(\xi)} A=\xi^{m} \frac{\partial A_{i j}}{\partial x^{m}}+A_{m j} \frac{\partial \xi^{m}}{\partial x^{i}}+A_{i m} \frac{\partial \xi^{m}}{\partial x^{j}}-\xi^{m} \frac{\partial A_{i j}}{\partial x^{m}}=0 \\
\nabla_{m(\xi)} A=A_{m j} \frac{\partial \xi^{m}}{\partial x^{i}}+A_{i m} \frac{\partial \xi^{m}}{\partial x^{j}}=0, \\
\text { ie } \quad A_{m j} \frac{\partial \xi^{a n}}{\partial x^{i}}+A_{i m} \frac{\partial \xi^{a m}}{\partial x^{j}}=0, \tag{7}
\end{gather*}
$$

and then after simplifying, we get

$$
\begin{align*}
& A_{11}(0)+A_{21}(0)+A_{31}(0)+A_{41}(0)+A_{11}(0)+A_{21}(0)+A_{31}(0)+A_{41}(0)=0 \\
& A_{12}=A_{21}=A_{13}=A_{31}=A_{22}=A_{23}=A_{32}=A_{24}=A_{42}=A_{33}=A_{34}=A_{43}=0 \tag{8}
\end{align*}
$$

Thus the bimetrically spherically symmetric covariant tensor $A_{i j}$ have the following non-vanishing components: $A_{11}, A_{14}, A_{41}, A_{44}$. They can be treated as functions of $r$ and $t$ to get Takeno's analogy

## MIXED TENSOR

In view of definition (1), mixed tensor $A_{j}^{i}$ of order-2 is bimetrically spherically symmetric if

$$
\begin{align*}
& \nabla_{m(\xi)} A_{j}^{i}=\left(g_{m}-f_{m}\right)_{(\xi)} A_{j}^{i}=0 \\
& \xi^{m} \frac{\partial A_{j}^{i}}{\partial x^{m}}+A_{m}^{i} \frac{\partial \xi^{m}}{\partial x^{j}}-A_{j}^{m} \frac{\partial \xi^{i}}{\partial x^{m}}-\xi^{m} \frac{\partial A_{j}^{i}}{\partial x^{m}}=0 \\
& A_{m}^{i} \frac{\partial \xi^{m}}{\partial x^{j}}-A_{j}^{m} \frac{\partial \xi^{i}}{\partial x^{m}}=0, \\
& \text { ie } \quad A_{m}^{i} \frac{\partial \xi^{a m}}{\partial x^{j}}-A_{j}^{m} \frac{\partial \xi^{a i}}{\partial x^{m}}=0,  \tag{9}\\
& \quad A_{1}^{1}(0)+A_{2}^{1}(0)+A_{3}^{1}(0)+A_{4}^{1}(0)-A_{1}^{1}(0)-A_{1}^{2}(0)-A_{1}^{3}(0)-A_{1}^{4}(0)=0  \tag{10}\\
& A_{1}^{2}(0)+A_{2}^{1}(0)+A_{3}^{1}(0)+A_{4}^{1}(0)+A_{2}^{2}(0)+A_{3}^{2}(0)+A_{4}^{2}(0)+A_{2}^{3}(0)  \tag{11}\\
& +A_{3}^{3}(0)+A_{4}^{3}(0)+A_{1}^{4}(0)+A_{2}^{4}(0)+A_{3}^{4}(0)+A_{4}^{4}=0
\end{align*}
$$

## CONCLUSION.

The following main themes have been sorted out.

1. The bimetric spherically symmetric scalar $A$ can be chosen as a function of $r$ and $t$ to mateh with Takeno (1966).
2. The bimetrically spherically symmetry vector $A_{i}$ assume the form
$\mathrm{A}_{\mathrm{i}}=\left(\mathrm{A}_{1}, 0,0, \mathrm{~A}_{4}\right)$, where $\mathrm{A}_{1}$ and $\mathrm{A}_{4}$ are the functions of (r, $\left.\theta, \phi, \mathrm{t}\right)$ To reclaim the results of Takeno (1966) one has to take, $\mathrm{A}_{1}$ and $\mathrm{A}_{4}$ independent of $\theta$ and $\phi$.
3. The bimetrically spherically symmetric covariant tensor $A_{i j}$ have the following non-vanishing components:
$A_{11}, A_{14}, A_{41}, A_{44}$. They can be treated as functions of $r$ and $t$ to get Takeno's analogy.
4. The bimetrically spherically symmetric mixed tensor $A_{i}^{j}$ have the following non-vanishing components:
$A_{1}^{1}, A_{1}^{4}, A_{4}^{1}, A_{4}^{4}$. They can be treated as functions of $r$ and $t$ to get Takeno's analogy.
Our investigation leads to the generalization of the Takeno's findings in the sense that the components under consideration comes out to be the functions of $r, \theta, \phi, t$. Also the major departure is with regard to the nonvanishing components as illustrated in the following table.

|  | Non-vanishing $A_{i}$ | Non-vanishing $A_{i j}$ |
| :--- | :--- | :--- |
| Takeno | $A_{1}, A_{4}$ | $A_{11}, A_{22}, A_{33}, A_{44}, A_{14}, A_{41}$ |
| Ours | $A_{1}, A_{4}$ | $A_{11}, A_{14}, A_{41}, A_{44}$ |

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